

Finite Element Modeling of Crustal Deformation in the North America–Caribbean Plate Boundary Zone

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velocity closure requirements predict only 1 ? mm/yr NA-Ca relative motion at [he. Cayman sprinting axis. *Deng and Sykes [1995]* used a least squares method to compute an NA-Ca Euler pole from NA-Ca PBZ earthquake slip vectors and then calculated an angular rotation rate such that it produces 20 mm/yr relative motion at the Cayman spreading center. They proposed that the mean rate derived from Cayman Trough spreading underestimates the full plate motion because of past or current slip along presumed Cayman fracture zones, and that motion directions based on Cayman transform faults may be unreliable if deformation has modified their azimuths [*Sykes et al., 1982*]. Thus, NUVEL-1 accounts for no distributed deformation in the.

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Mann et al., 1991]. Seismicity along the fault (Figure 2) clearly indicates pure left-lateral slip. We infer from the simplicity of Swan Islands Fault structures and seismicity that the fault is the plate boundary and that motion here is not distributed across a wide zone as it is further east. The Swan Islands Fault terminates eastward at the southern end of the Cayman spreading center, which has normal faulting earthquakes (Figure 2) and generates the magnetic anomalies of the Cayman Trough [*Rosencrantz et al.*, 1988]. The northern end of the Cayman spreading axis is

underthrusting of Caribbean lithosphere beneath terranes including eastern Hispaniola and Puerto Rico [Byrne *et al.*, 1985; Armstrong and Pennington, 1990; Masson and Scanlon 1991], South of Puerto Rico there is evidence that Muertos Trough convergence decreases eastwards [Masson and Scanlon, 1991] with strike-slip motion [P. Jansma, *personal communication*, 1995] near its intersection with the

where in the finite element array. A Monte Carlo technique is used to interpret the sensitivity and uncertainties in the model due to complexities in the data. For our application to the northern Caribbean region, we solve for nodal displacements on an elastic, 2-D spherical shell with a plane stress condition ($\sigma_3 = 0$). We use eight node, biquadratic, isoparametric elements.

The crustal layer is modeled as a two dimensional elastic plate, cut by active faults. Faults are surfaces in which 1) only the slip direction is constrained; 2) both the slip direction and rate are known; or 3) both the direction and rate are unconstrained. Fault slip data are implemented as described by *Melosh and Raefsky [1981]*, and introduce torque-free dislocations in the continuum. Fault orientations are prescribed as free-shear surfaces [*Melosh and Williams, 1989*].

Far field displacements are imposed as boundary conditions. For a spherical geometry, relative plate motions imposed at the mesh boundary are calculated based on a given boundary node's angular distance from the Euler pole. Likewise, the overall orientation of the finite element mesh boundaries are controlled by the Euler pole location since we choose to have the boundaries fall on either great or small circles relative to the rotation pole. Thus, when we compare models based on different poles, the locations of the mesh boundaries reflect differences in pole location.

The most significant difference between this approach and the rigid

on a spherical shell is:

$$\varepsilon = \begin{Bmatrix} \varepsilon_\theta \\ \varepsilon_\phi \\ 2\varepsilon_{\theta\phi} \end{Bmatrix} = \frac{1}{R} \left\{ \begin{bmatrix} \frac{\partial}{\partial \theta} & 0 \\ 0 & \frac{1}{\sin \theta} \frac{\partial}{\partial \phi} \\ \frac{1}{\sin \theta} \frac{\partial}{\partial \phi} & \frac{\partial}{\partial \theta} \end{bmatrix} + \begin{bmatrix} 0 & 0 \\ \cot \theta & 0 \\ 0 & -\cot \theta \end{bmatrix} \right\} \begin{Bmatrix} u_\theta \\ u_\phi \end{Bmatrix} = Su + SN\zeta = B\zeta \quad (2)$$

where, (θ, ϕ) are the colatitude and longitude, respectively, and R is the radius of the shell (radius of the earth). For an elastic body in equilibrium, we use a variational principle to solve for the minimum in the potential energy with respect to small nodal displacements. For a given element, e , the elastic strain energy is

$$U^e = \frac{1}{2} \int_{\Omega^e} \varepsilon^T \sigma d\Omega \quad (3)$$

and using Hooke's law (1),

$$U^e = \frac{1}{2} \int_{\Omega^e} \varepsilon^T D \varepsilon d\Omega = \frac{1}{2} \int_{\Omega^e} \zeta^T B^e T D^e B^e \zeta d\Omega \quad (4)$$

where Ω^e is the element volume. For a boundary traction vector t applied on the boundary of the element, the potential energy of the applied boundary stress is

$$U^e = - \int_S u^T t dS = - \int_S [N \zeta^e]^T t dS \quad (5)$$

where S is the surface of the element. Summing (4) and (5) and minimizing the total energy of the body under load with respect to small variations in ζ , we get the standard finite element equation:

$$K\zeta = f \quad (6)$$

where

$$K^e = \int_{\Omega^e} B^e T D^e B^e d\Omega \quad (7)$$

defines the element stiffness matrix and

$$f^e = \int_S [N^e]^T t dS \quad (8)$$

defines the element force vector. The global stiffness matrix K is assembled from K^e and is solved for by Gaussian integration. Displacement boundary conditions and fault slip data (split nodes) enter into (8) which reduces the rank to the number of degrees of freedom. A sufficient condition for K to be inverted is that the translation and rotation for each element are constrained so the solution $\zeta = K^{-1}f$

exists.

We calculate uncertainties in the model predictions via a Monte Carlo method to construct the probability distribution for the deformation everywhere. This is computationally efficient since the kinematic data enter the model in the force vector \mathbf{f} , so the stiffness matrix, \mathbf{K} , is inverted only once. Histograms are constructed from typically

Discussion

The NUVEL-1-based model (where mesh orientation and far-field boundary rates are based on the NUVEL-1 model of *DeMets et al.*, 1990; see Figure 3) is the only model which clearly does not fit the observed left-lateral motion along the Walton-Plantain Garden-Enriquillo fault system [*Mann et al.*, 1995] and the observed compression at the Jamaica restraining bend [*Burke et al.*, 1980; Fig. 2 focal mechanism]. This is not unexpected given that its far field NA-Ca rate is less than the observed Cayman spreading rate. Despite this, we note how similar the NUVEL-1 derived motion patterns are in the eastern

Septentrional to the eastern Enriquillo fault, in central Hispaniola gives slow thrust motion from 1-4 mm/yr along its length, consistent with field observations [Mann *et al.*, 1995].

The C&M and D&S models differ most significantly in the easternmost PBZ between Puerto Rico and the Aneгада Passage fault system. Morphology, bathymetry, and seismicity indicate that the Aneгада Passage fault system is an actively extending graben [Murphy and McCann, 1979; Frankel *et al.*, 1980; Jany *et al.*, 1990; Masson and Scanlon, 1991; Speed and Larue, 1991]. The Aneгада Passage extension direction is not well constrained by earthquake focal mechanisms, right-lateral component along the fault system has been proposed

GPS site FRAN has ESE directed motion. If we insist on a pure strike-slip motion on the central portion of the Septentrional fault, we can find a solution (Figure 8) in which the crustal motion at FRAN is ESE-directed, and a normal component of motion is transferred to the fault system off the NE coast of Hispaniola. This result indicates that there is both vertical and horizontal strain partitioning in the restraining bend here: the thrust motions observed by *Russo and Villaseñor* [1995] (see Fig. 2) occur on a deeper crustal fault or faults beneath NE Hispaniola, and the surficial

cate that NA-Ca relative motion is partitioned as displacement along faults and as continuum deformation (including rotation) throughout the NA-Ca plate boundary zone. Our results are in good agreement with preliminary GPS results [Farina *et al.*, 1995], if we constrain the central Septentrional fault to left-lateral strike-slip motion.

Acknowledgements

We thank Pam Jansma for generously sending us unpublished data and for helpful discussion. Bob Speed read and helped correct flaws in an early version of this work. We thank Eric Calais, Chuck DeMets, and Paul Mann, for critical reviews of the manuscript. Bill McCann offered useful advice. We thank Guilhem Barruol for use of computer facilities and logistical help. A research associateship at the Department of Terrestrial Magnetism and NSF Grant EAR 93-16-1 S") to P. Silveranti R. Russo supported R. Russo during this work. Support for P. Lundgren was provided by the Jet Propulsion Laboratory, California Institute of Technology, under contract to the National Aeronautics and Space Administration.

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Figure 4b. *Deng and Sykes [1995]* based model: fault rate solution (see Figure 3b for explanation).

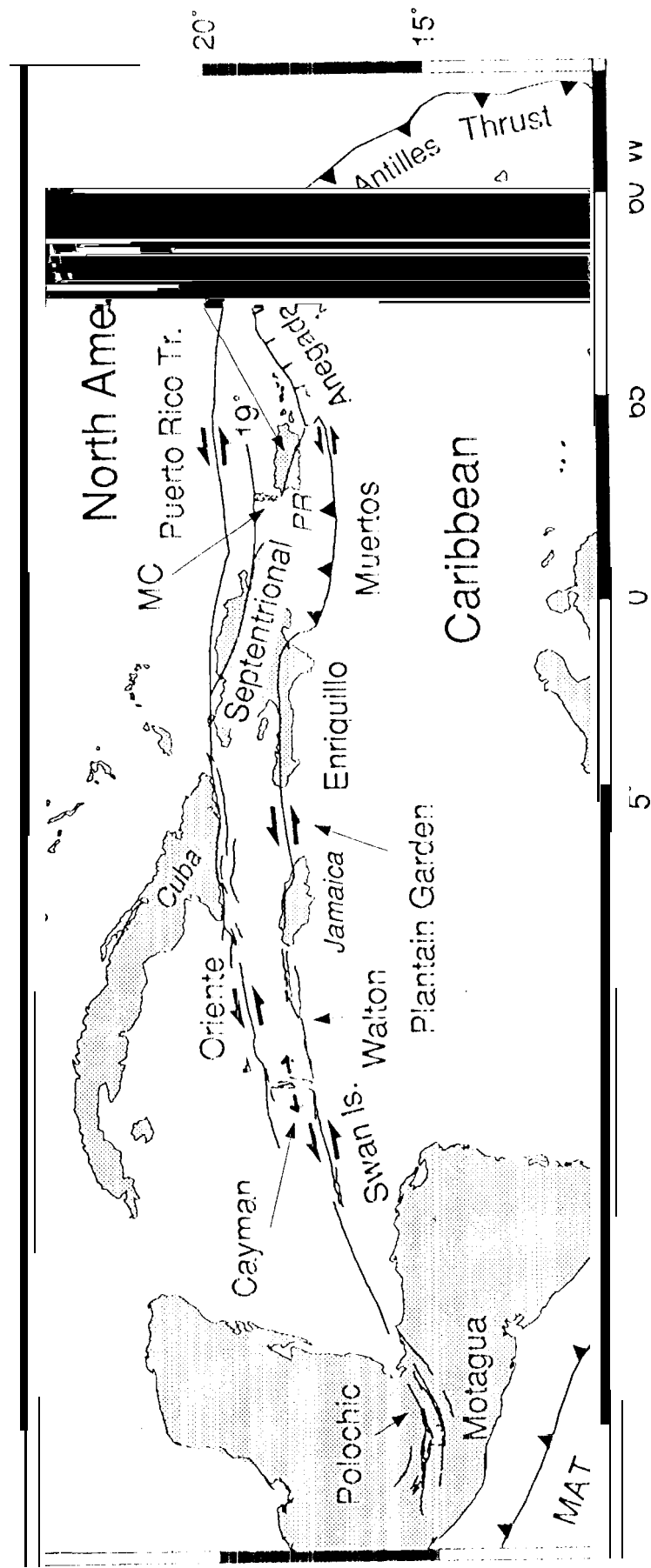
Figure 5a. *Calais and Mercier de Lépinay [1993]* based model: calculated motions at the center of each element (see Figure 3a for explanation)

Figure 5b. *Calais and Mercier de Lépinay [1993]* based model: fault rate solution (see Figure 3b for explanation).

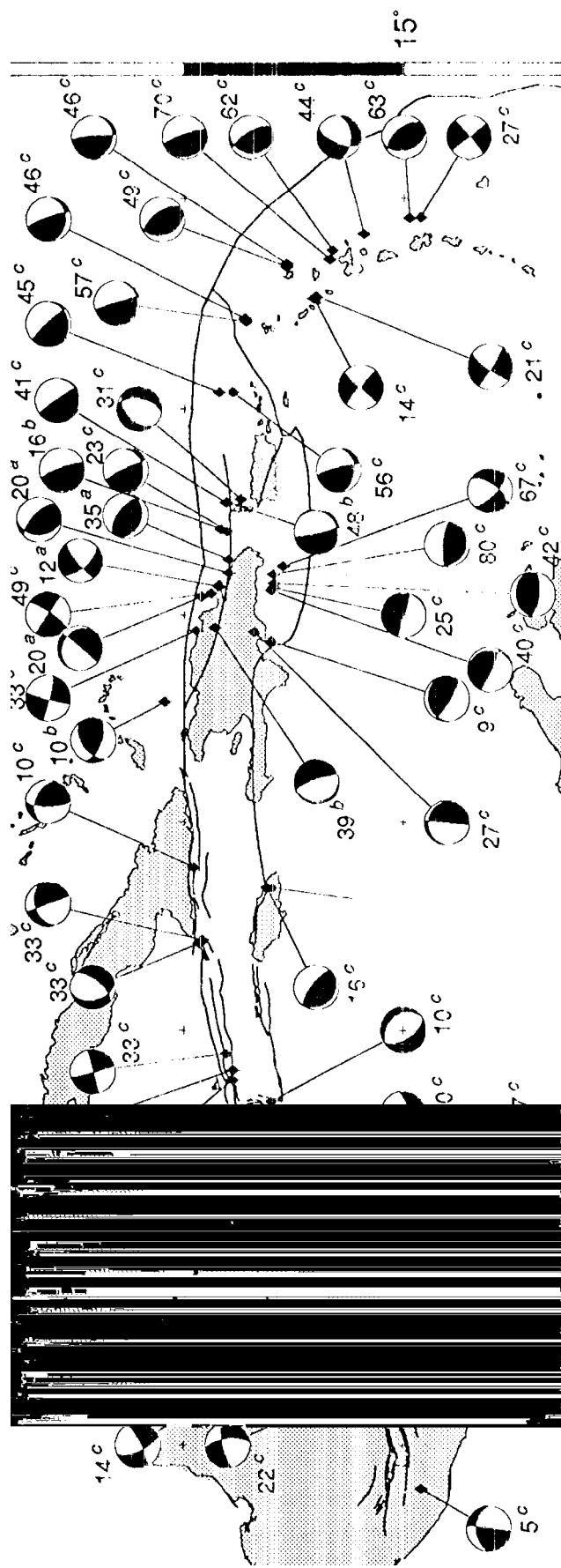
Figure 6. Preliminary GPS site velocity vectors and their 95% confidence ellipses with respect to site TURK in the Bahamas on stable North America plate for two occupations in 1986 and 1994 from *Farina et al. [1995]*.

Figure 7 *Calais and Mercier de Lépinay [1993]* based model (inset of Puerto Rico area): calculated motions at the center of each element (see Figure 3a for explanation). This is the same solution as Figure 5a but with the North America plate held fixed so that the continuum velocity vectors can be compared to Figure 6.

Figure 8 *Calais and Mercier de Lépinay [1993]* based model (inset of Puerto Rico area): (a) calculated motions at the center of each element (see Figure 3a for explanation). This is the same as Figure 7 but with the added constraint that the central Septentrional fault be strike-slip. This produces continuum motions more consistent with the GPS results shown in Figure 6. (b) fault rate solution (see Figure 3b for explanation).



Figure



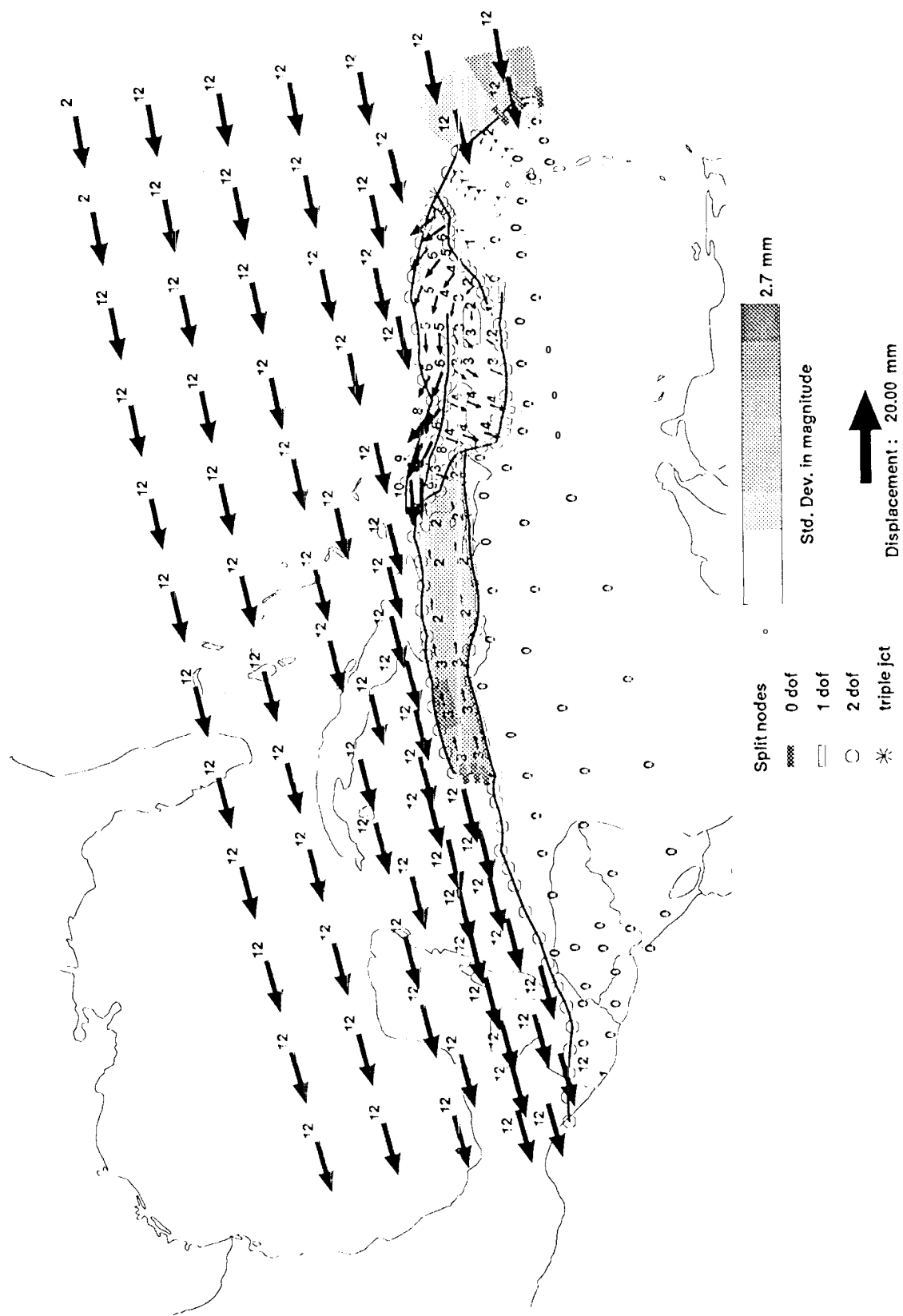


Figure 3a

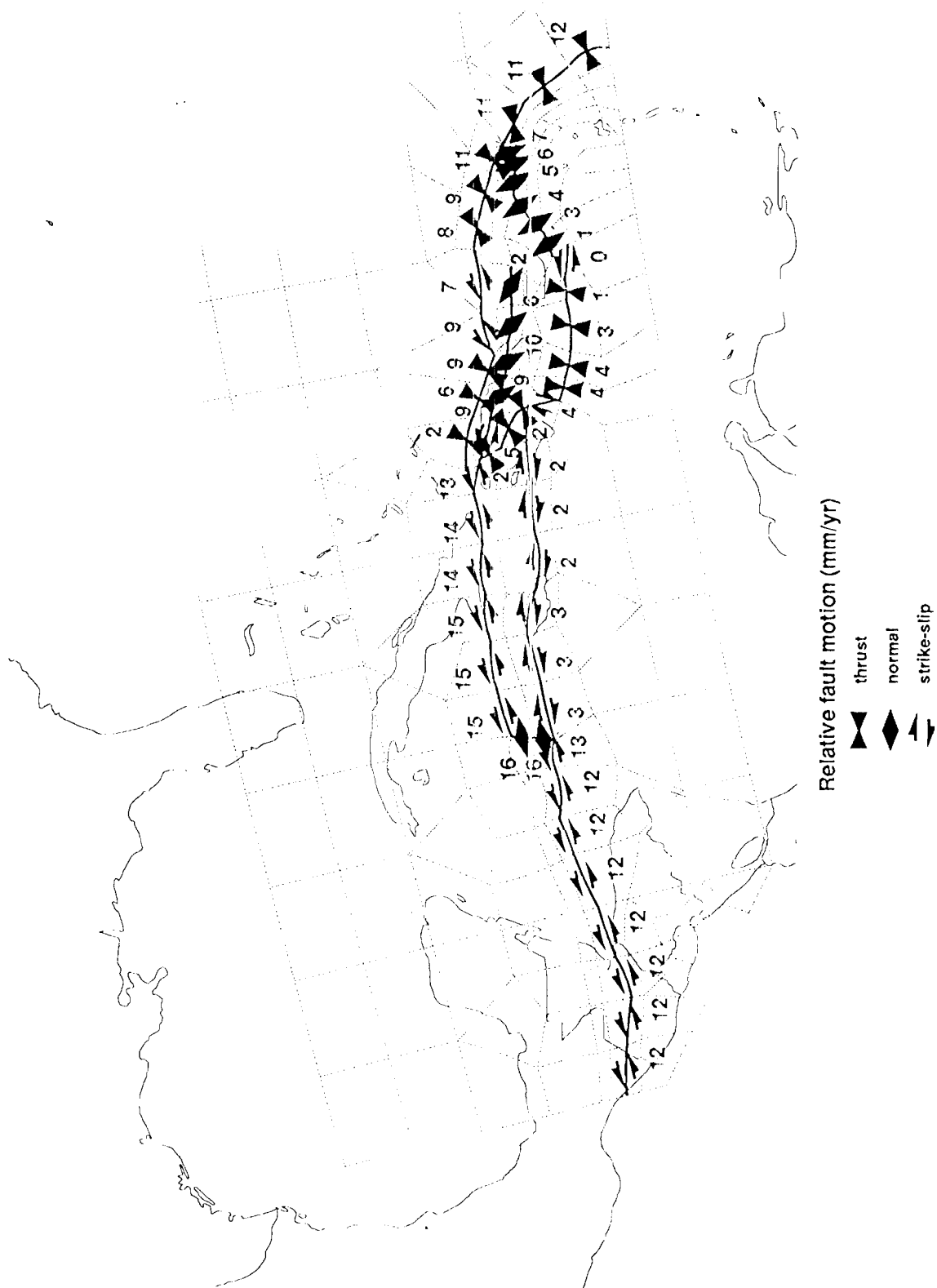


Figure 3b

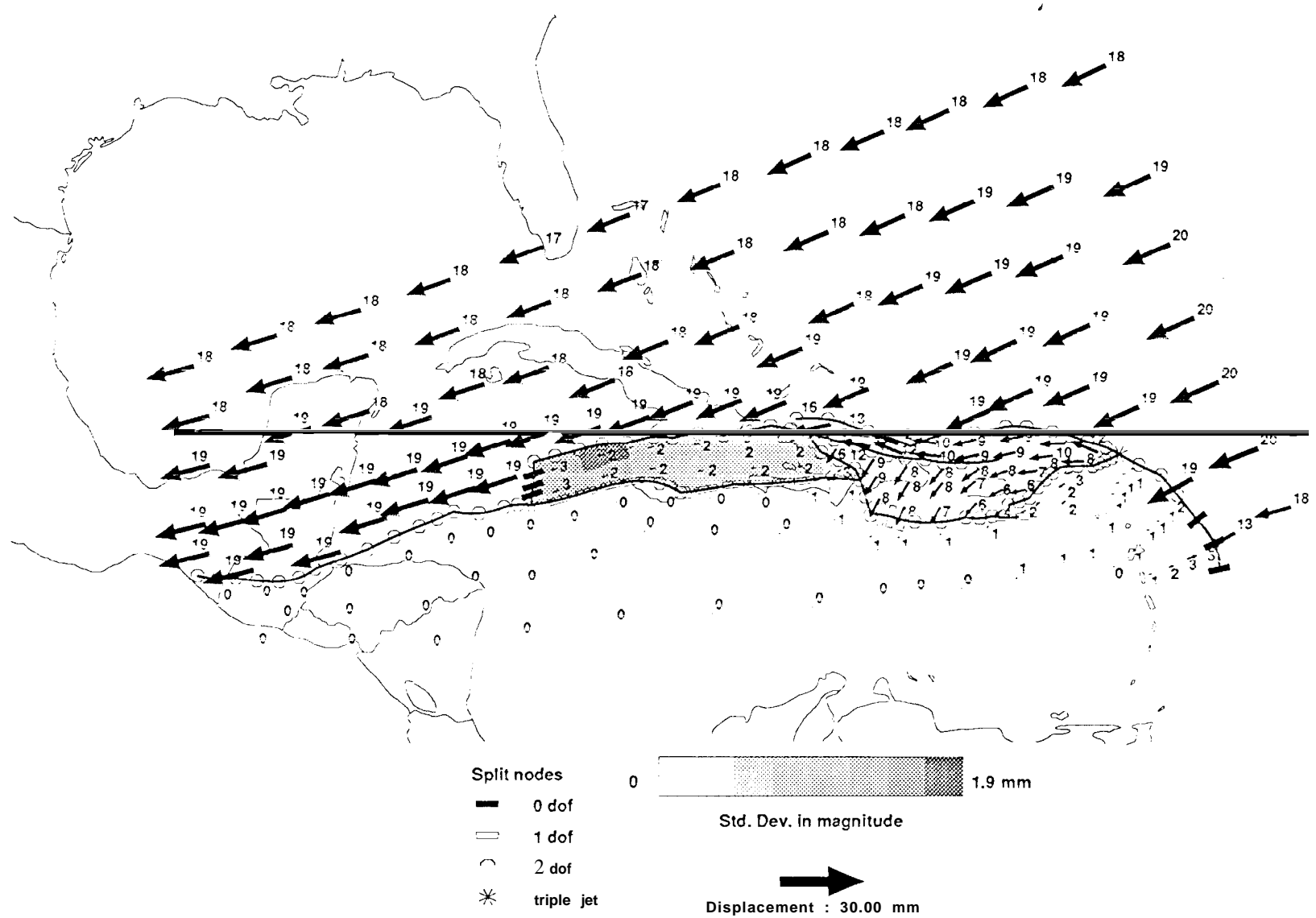


Figure 4a



Figure 4b

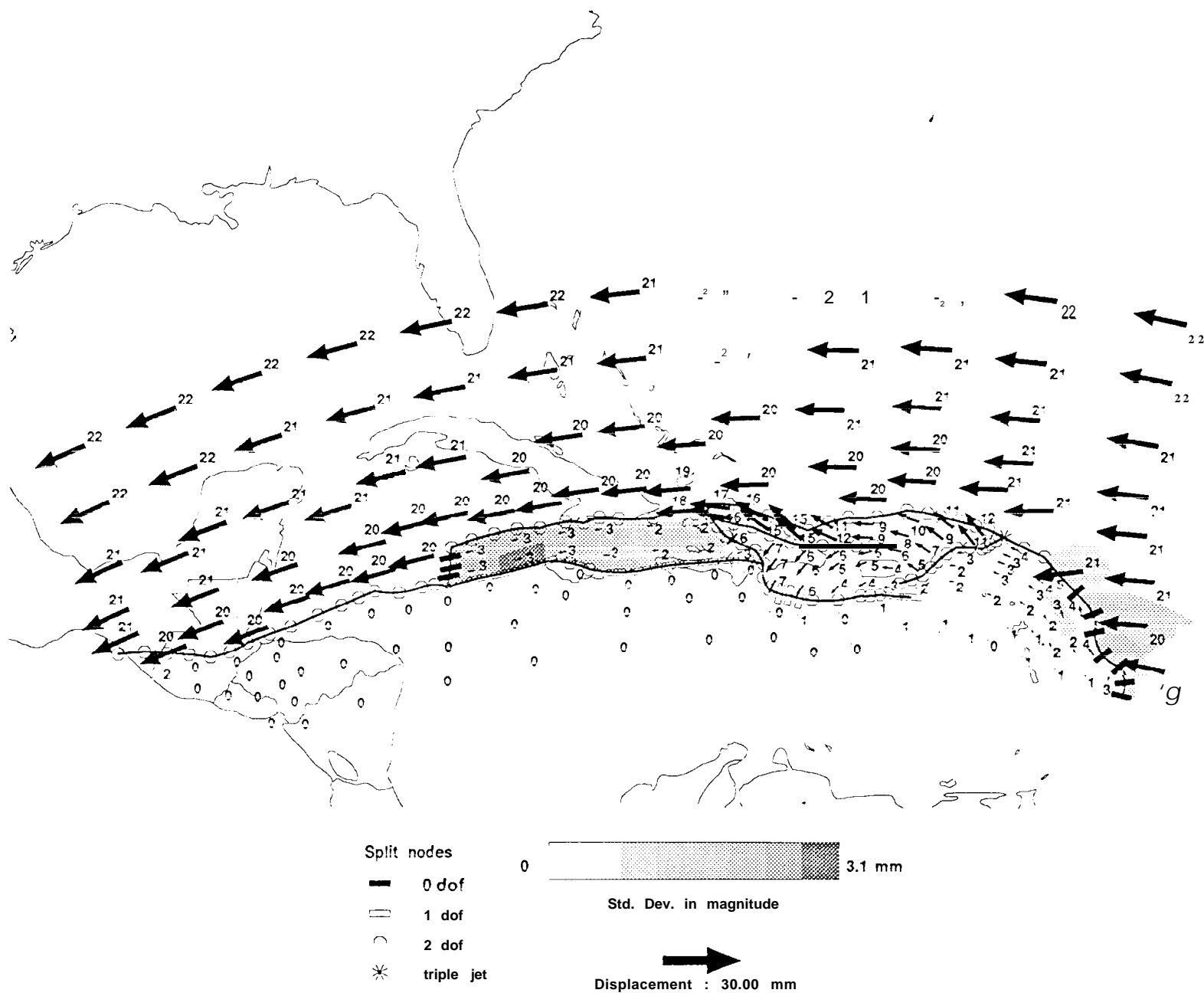


Figure 5a

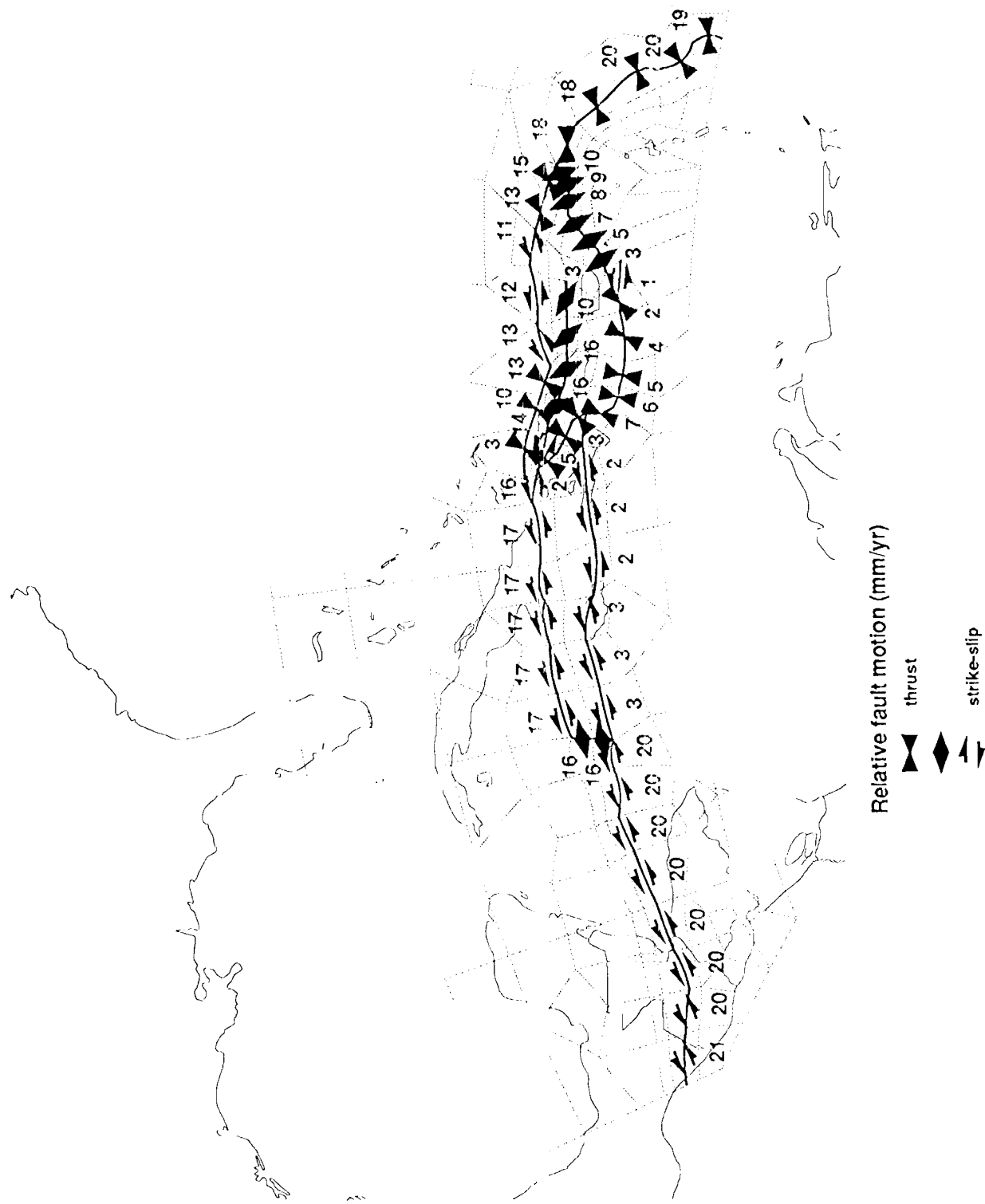


Figure 5b

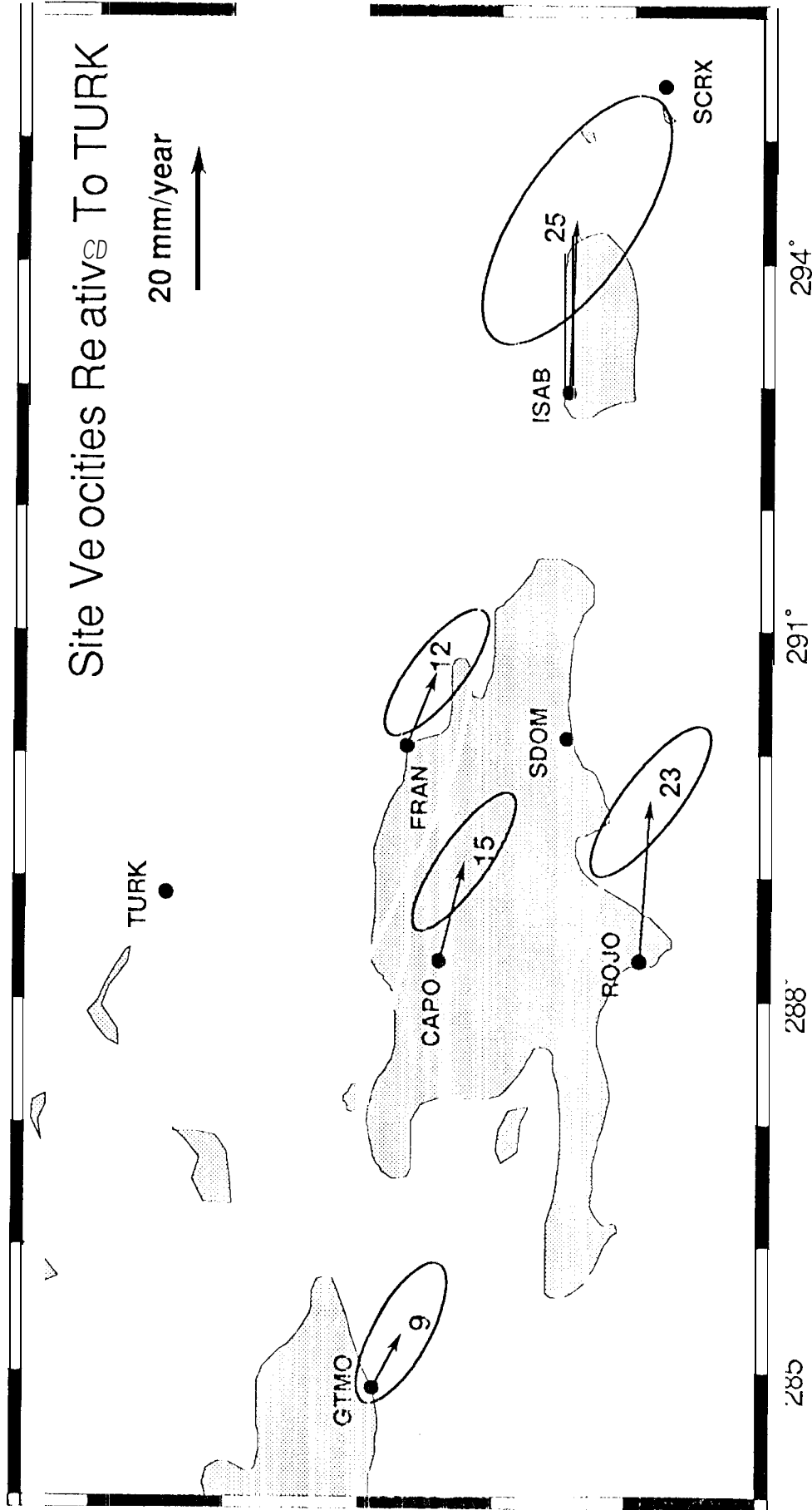


Figure 6

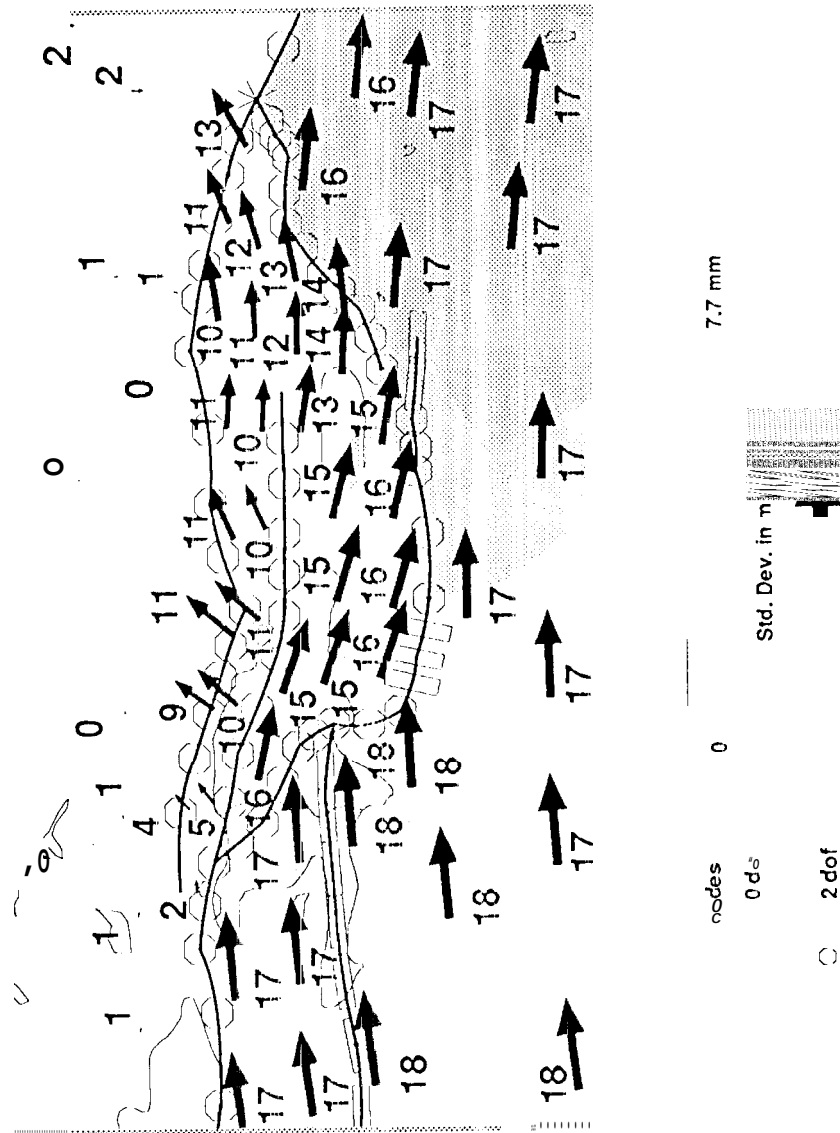


Fig 7

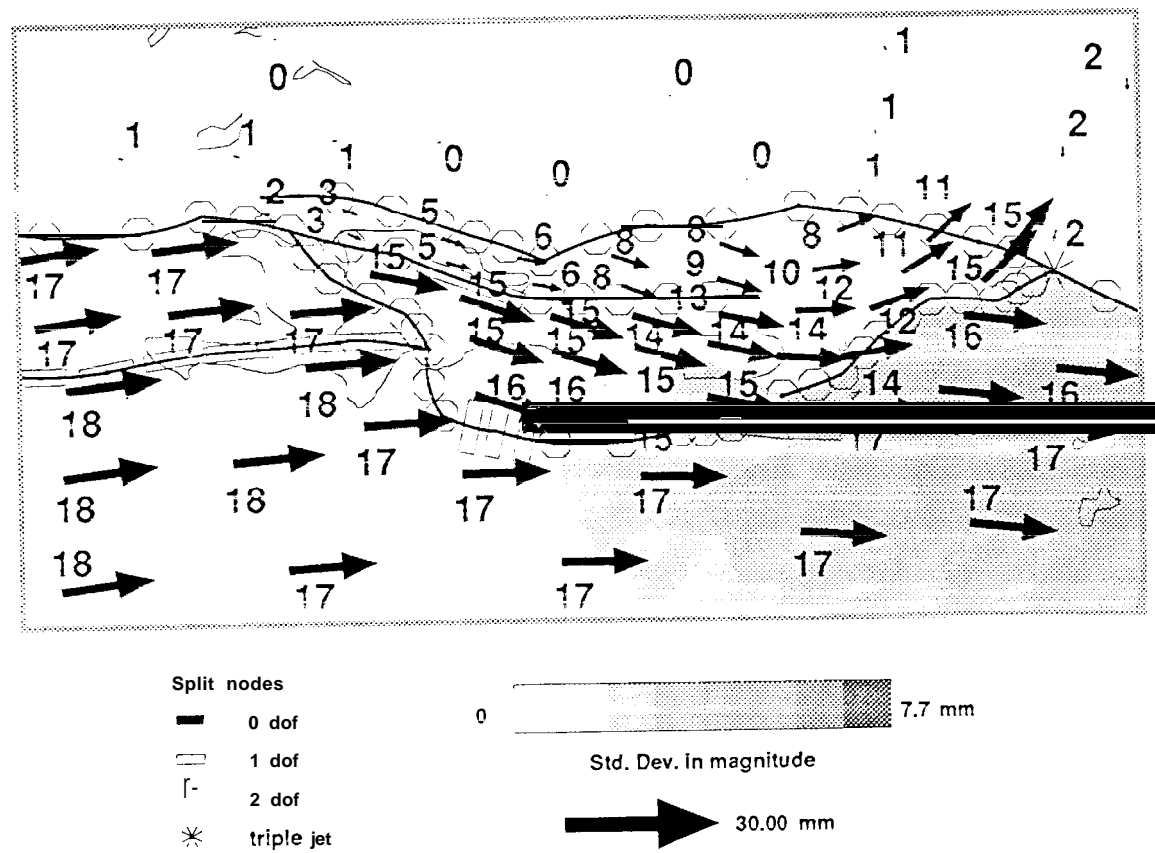


Figure 8a

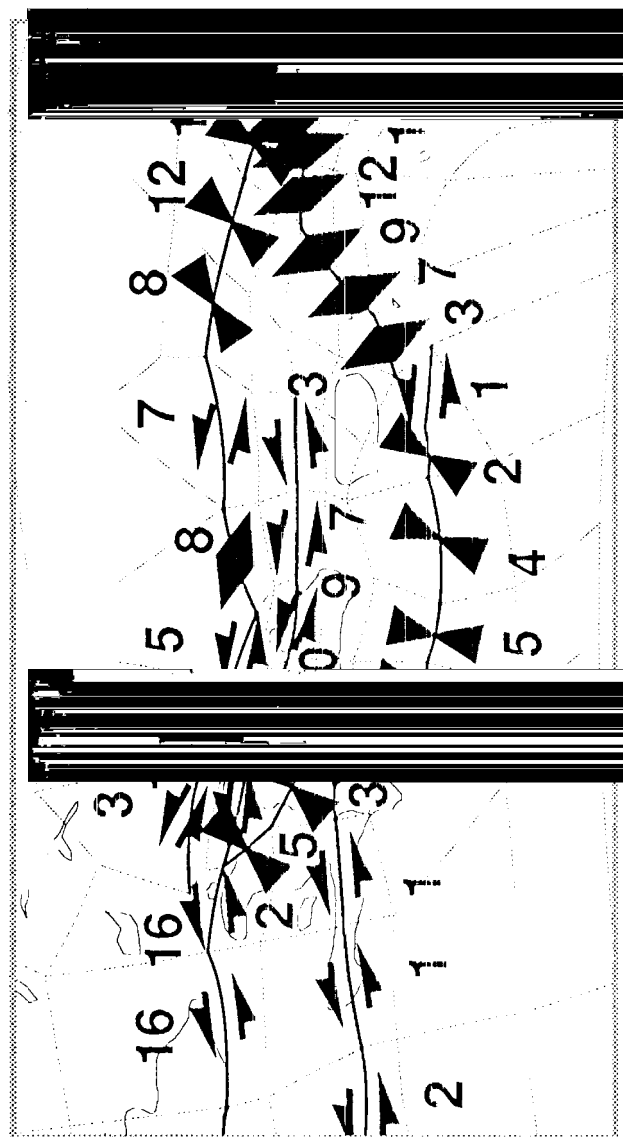


Fig. e 8b